

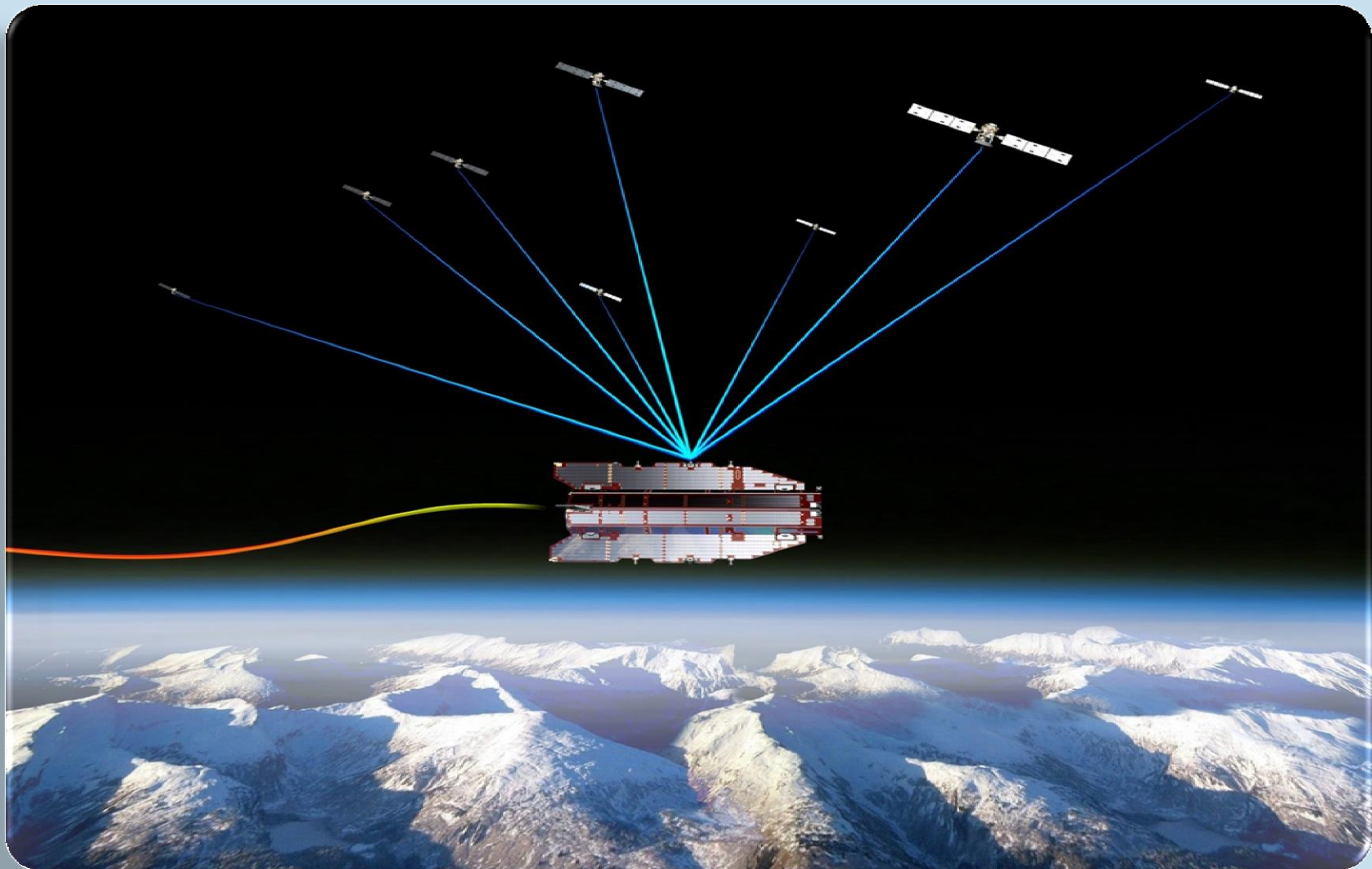
# A New Approach for Pure Kinematical and Reduced Kinematical Determination of a LEO Orbit based on GNSS Observations

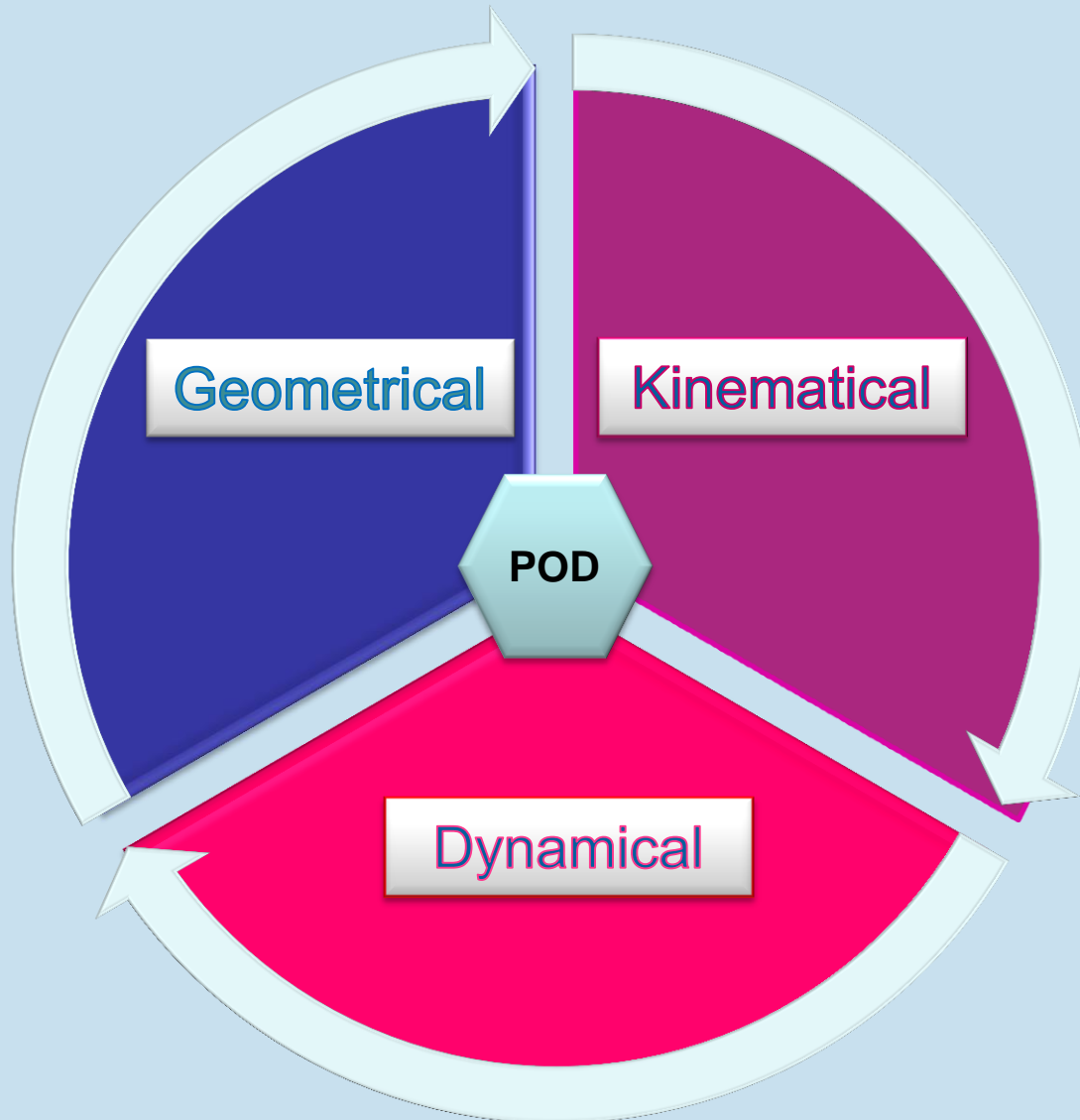
Akbar Shabanloui

Session: Gravity 2.2

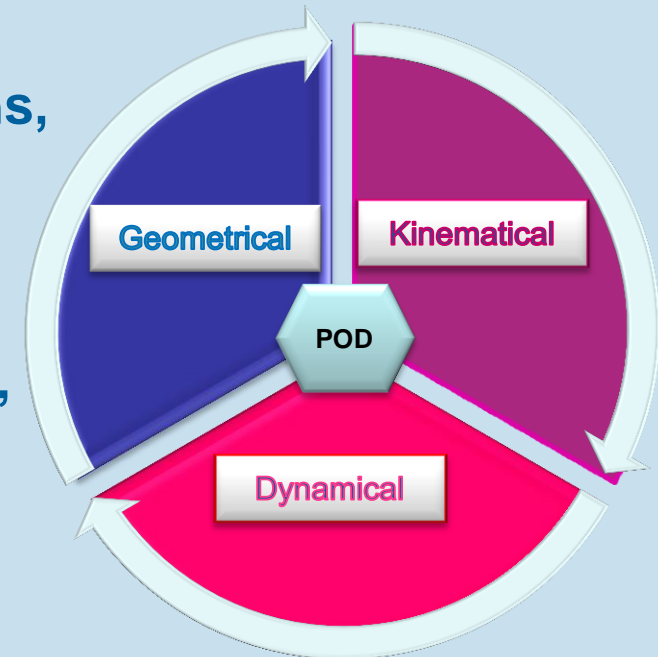
IAG 2009

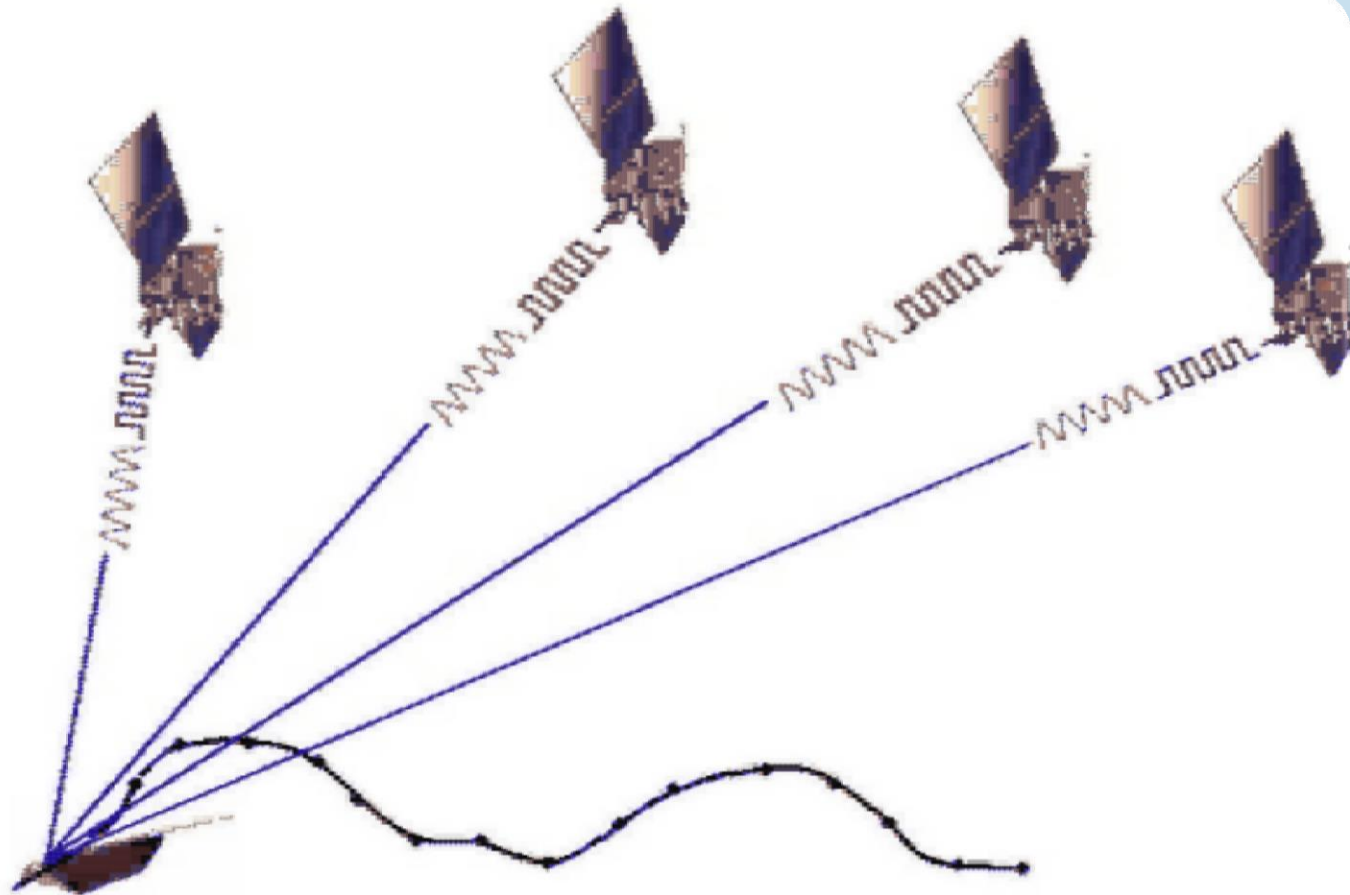
31th August 2009, Buenos Aires





- **Geometrical POD** : point-wise, positions **!= Kinematical POD**
- **Kinematical POD** : continous, positions, velocities and accelerations
- **Dynamical POD** : continous, positions, velocities and accelerations based on force function information





GPOD

KPOD

RKPOD

DPOD

Method

Method

Simulation

Simulation

Real data

Real data">

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graph TD; GPOD[GPOD]; KPOD[KPOD]; RKPOD[RKPOD]; DPOD[DPOD]; KPOD --- KPOD_Method[Method]; KPOD --- KPOD_Simulation[Simulation]; KPOD --- KPOD_RealData[Real data]; RKPOD --- RKPOD_Method[Method]; RKPOD --- RKPOD_Simulation[Simulation]; RKPOD --- RKPOD_RealData[Real data];
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GPOD

KPOD

RKPOD

DPOD

Method

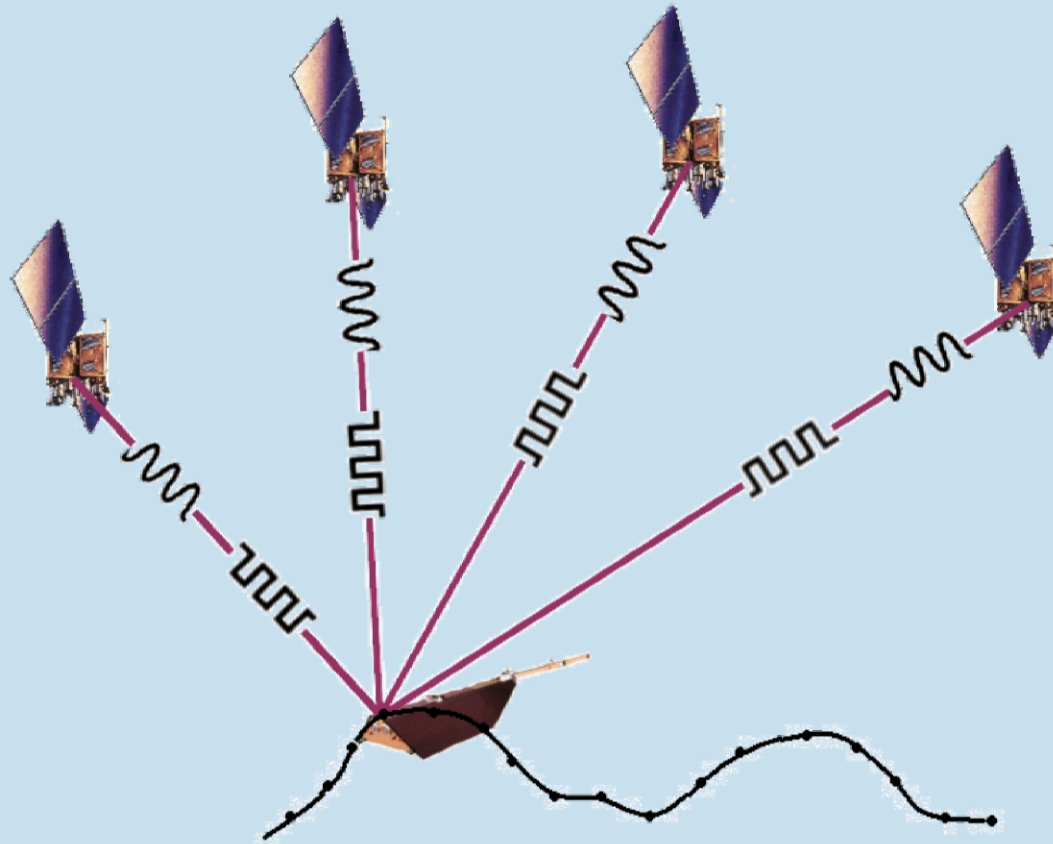
Method

Simulation

Simulation

Real data

Real data



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Equation of motion w.r.t self-adjoint differential operator:

$$L(\mathbf{r}(t)) = \mathbf{a}(t; \mathbf{r}, \dot{\mathbf{r}})$$

corresponding Fredholm's integral equation:

Reference motion    Integral kernel

$$\mathbf{r}(\tau) = \bar{\mathbf{r}}(\tau) - T^2 \int_{\tau'=0}^1 K(\tau, \tau') \mathbf{a}(\tau'; \mathbf{r}, \dot{\mathbf{r}}) d\tau'$$

Position
 $T = t_B - t_A$ 
 $\tau = \frac{t - t_A}{T}$ 
Force function

$$L(\bar{\mathbf{r}}(\tau)) = 0$$

A satellite short arc:

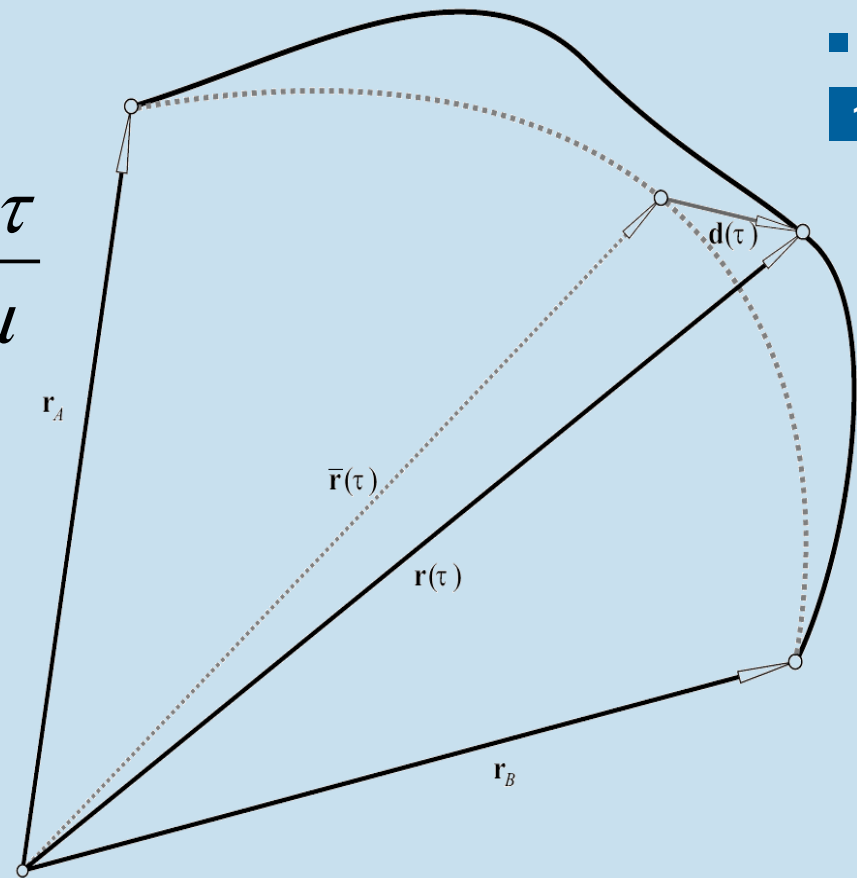
$$\mathbf{r}(\tau) = \bar{\mathbf{r}}(\tau) + \mathbf{d}(\tau) = \bar{\mathbf{r}}(\tau) + \sum_{\nu=1}^{\infty} \mathbf{d}_{\nu} \sin(\nu\pi\tau)$$

elliptical reference motion:

$$\bar{\mathbf{r}}(\tau) = \mathbf{r}_A \frac{\sin \mu(1 - \tau)}{\sin \mu} + \mathbf{r}_B \frac{\sin \mu\tau}{\sin \mu}$$

difference function:

$$\mathbf{d}(\tau) = \sum_{\nu=1}^{\infty} \mathbf{d}_{\nu} \sin(\nu\pi\tau)$$

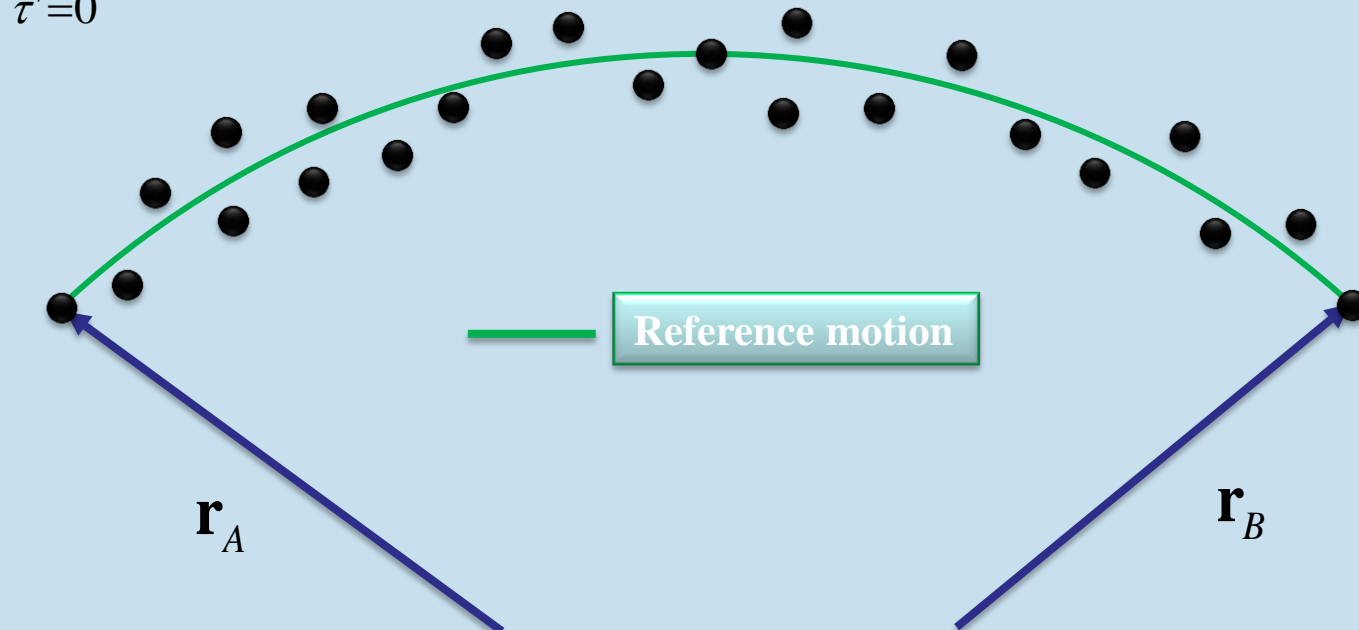


A satellite short arc can be represented:

$$\mathbf{r}(\tau) = \bar{\mathbf{r}}(\tau) + \mathbf{d}(\tau) = \bar{\mathbf{r}}(\tau) + \sum_{\nu=1}^{\infty} \mathbf{d}_{\nu} \sin(\nu\pi\tau)$$

$$\mathbf{r}(\tau) - \bar{\mathbf{r}}(\tau) = \mathbf{d}(\tau) = \sum_{\nu=1}^{\infty} \mathbf{d}_{\nu} \sin(\nu\pi\tau)$$

$$\mathbf{d}_{\nu} = 2 \int_{\tau'=0}^1 \mathbf{d}(\tau) \sin(\nu\pi\tau') d\tau'$$

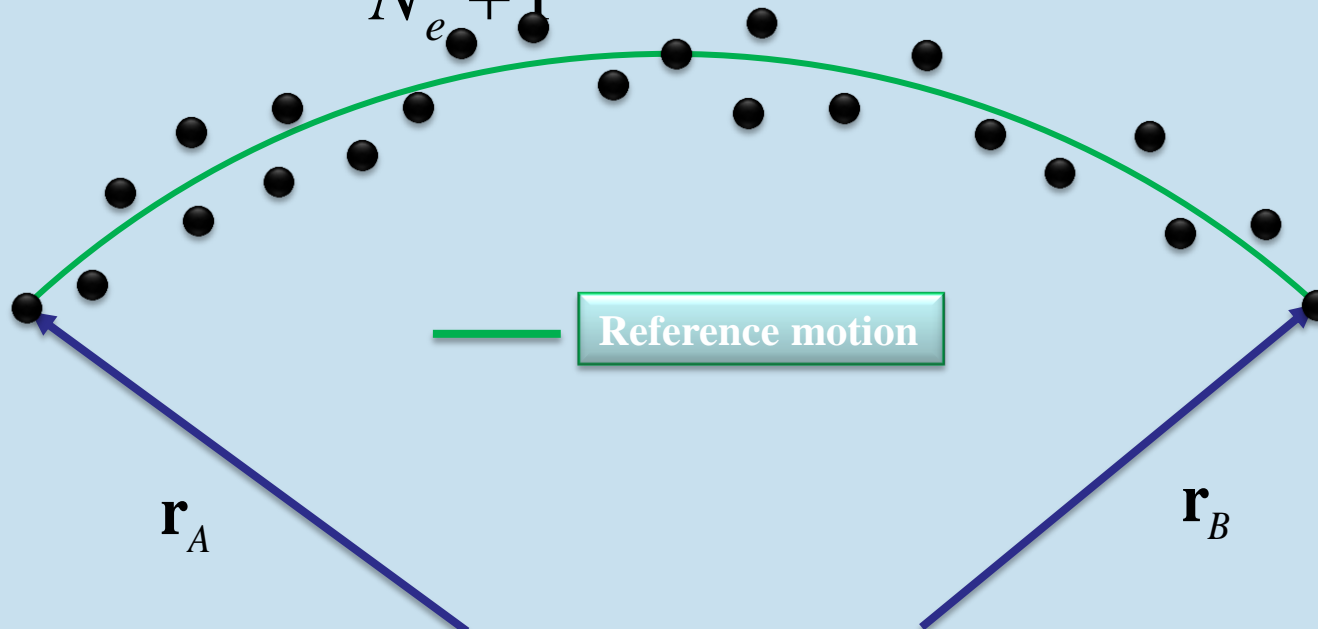


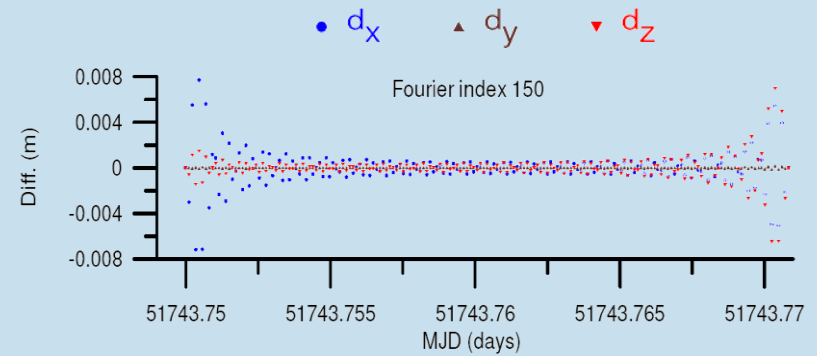
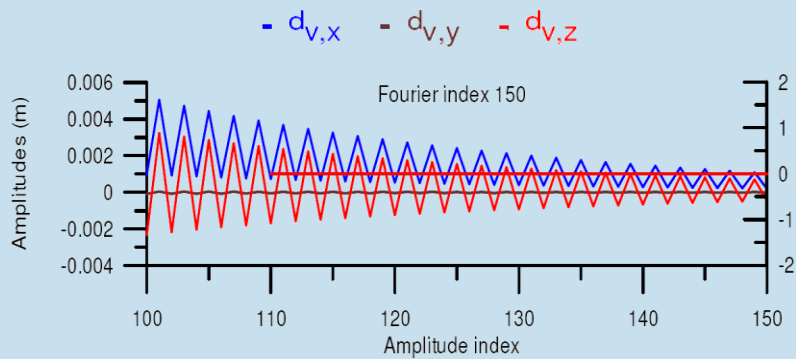
A satellite short arc can be represented:

$$\mathbf{r}(\tau) - \bar{\mathbf{r}}(\tau) = \mathbf{d}(\tau) = \sum_{\nu=1}^{\infty} \mathbf{d}_{\nu} \sin(\nu\pi\tau)$$

$$\mathbf{d}_{\nu} \approx \frac{2}{N_e + 1} \sum_{e=1}^{N_e} \mathbf{d}(\tau_e) \sin\left(\frac{\nu\pi e}{N_e + 1}\right)$$

$$\mathbf{d}(\tau_e) \quad \text{with} \quad \tau_e = \frac{e}{N_e + 1}, \quad e = 1, \dots, N_e \quad \tau_e \in ]0, 1[$$





ellipse mode

Amplitudes

Remainders

A satellite short arc can be represented:

$$\mathbf{r}(\tau) = \bar{\mathbf{r}}(\tau) + \sum_{\nu=1}^{\infty} \mathbf{d}_{\nu} \sin(\nu\pi\tau)$$

with Fourier amplitudes:

$$\mathbf{d}_{\nu} = 2 \int_{\tau'=0}^1 \mathbf{d}(\tau) \sin(\nu\pi\tau') d\tau'$$

Fourier series amplitudes:

$$\begin{aligned} \mathbf{d}_{\nu} = & \sum_{j=1}^J \frac{2(-1)^{j+1}}{(\nu\pi)^{2j+1}} [(-1)^{\nu} \mathbf{d}^{[2j]}(1) - \mathbf{d}^{[2j]}(0)] + \\ & + \beta \frac{2}{(\nu\pi)^{2J+1}} \int_{\tau'=0}^1 \mathbf{d}^{[2J+2]}(\tau') \sin(\nu\pi\tau') d\tau' \end{aligned}$$

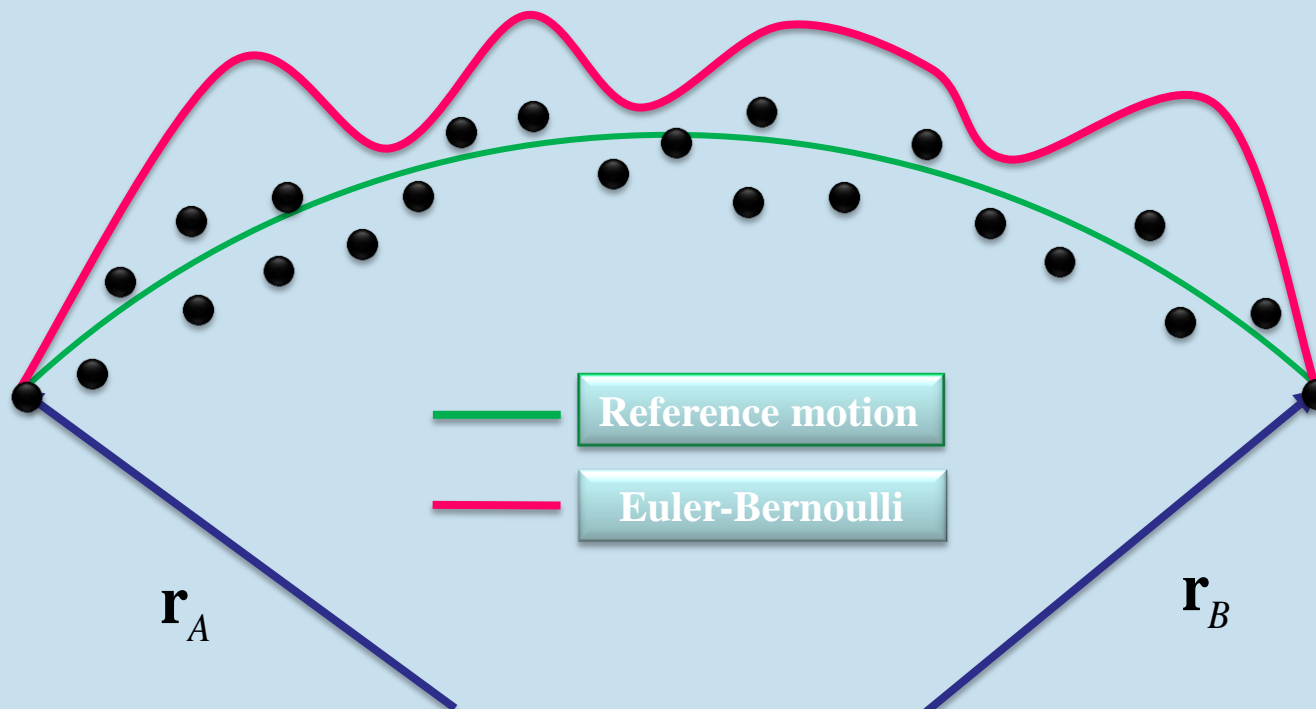
$$\begin{aligned}\mathbf{d}_F^\infty = \mathbf{d}(\tau) &= \sum_{\nu=1}^{\infty} \mathbf{d}_\nu \sin(\nu\pi\tau) = \\ &= \sum_{j=1}^{\infty} \mathbf{e}_{2j} E_{2j}(\tau) + \sum_{j=1}^{\infty} \mathbf{b}_{2j+1} B_{2j+1}(\tau) = \mathbf{d}_P^\infty\end{aligned}$$

A satellite short arc can be represented:

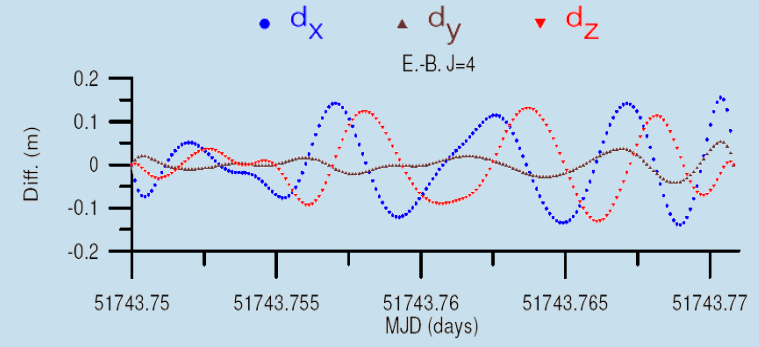
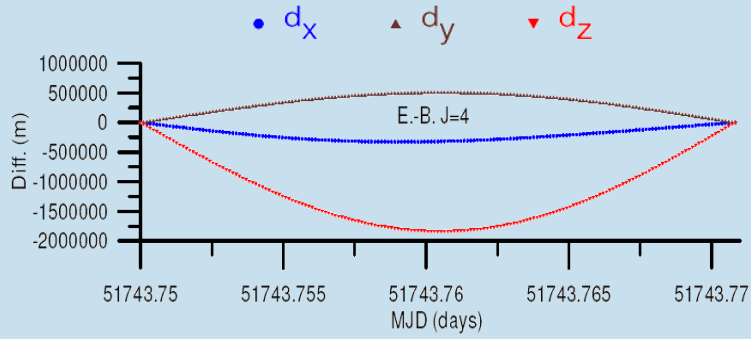
$$\sum_{\nu=1}^{\infty} \mathbf{d}_\nu \sin(\nu\pi\tau) = \sum_{j=1}^{\infty} \mathbf{e}_{2j} E_{2j}(\tau) + \sum_{j=1}^{\infty} \mathbf{b}_{2j+1} B_{2j+1}(\tau)$$

A satellite short arc can be represented with the Euler-Bernoulli term up to degree  $J$  as:

$$\mathbf{r}(\tau) - \bar{\mathbf{r}}(\tau) = \mathbf{d}(\tau) \approx \sum_{j=1}^J \mathbf{e}_{2j} E_{2j}(\tau) + \sum_{j=1}^J \mathbf{b}_{2j+1} B_{2j+1}(\tau)$$







ellipse mode

E.-B terms

Remainders

LEO orbit can be represented as:

$$\mathbf{r}(\tau) = \bar{\mathbf{r}}(\tau) + \mathbf{d}(\tau) = \bar{\mathbf{r}}(\tau) + \sum_{\nu=1}^n \mathbf{d}_{\nu} \sin(\nu\pi\tau)$$

**Gibbs effect!**

or

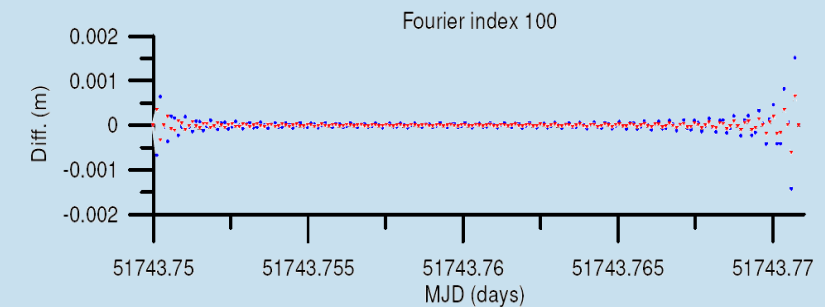
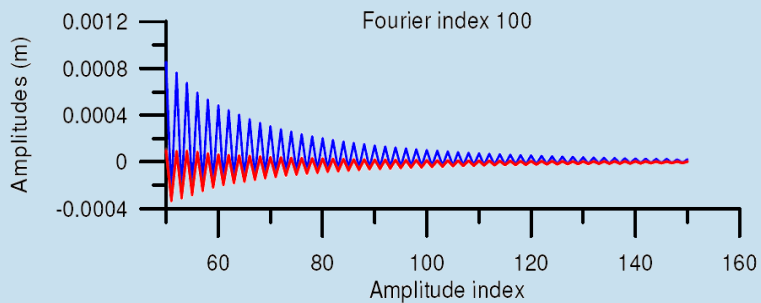
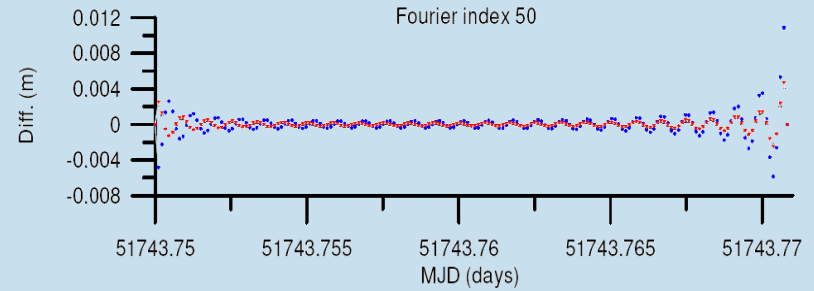
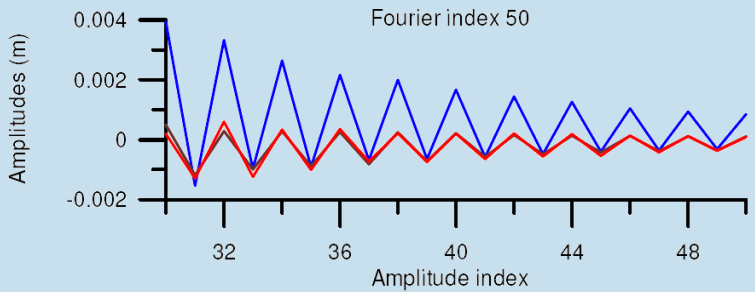
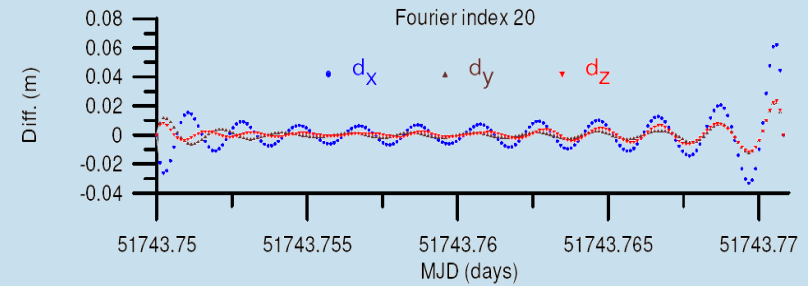
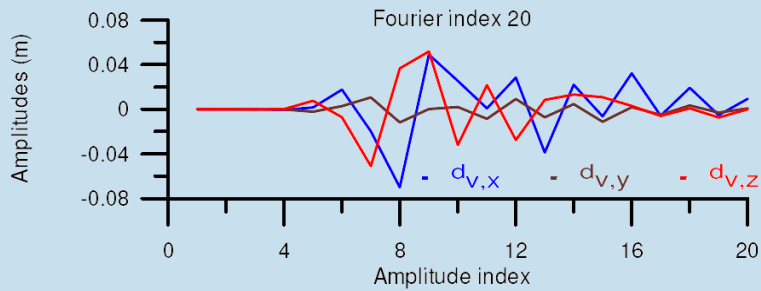
$$\mathbf{r}(\tau) = \bar{\mathbf{r}}(\tau) + \mathbf{d}(\tau) = \bar{\mathbf{r}}(\tau) + \sum_{j=1}^J \mathbf{e}_{2j} E_{2j}(\tau) + \sum_{j=1}^J \mathbf{b}_{2j+1} B_{2j+1}(\tau)$$

**Precision!**

**Solution?**

**fast  
convergence!**

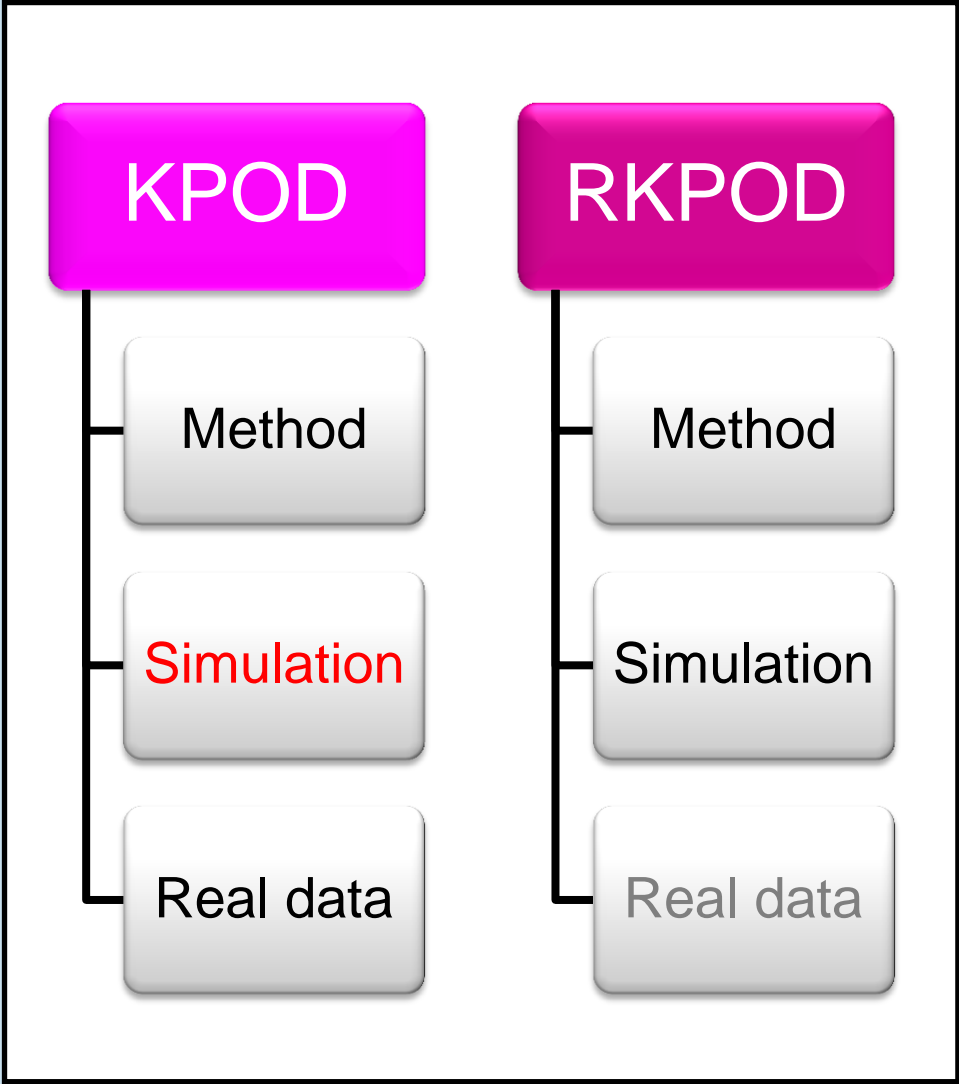
$$\mathbf{r}(\tau) = \bar{\mathbf{r}}(\tau) + \sum_{j=1}^J \mathbf{e}_{2j} E_{2j}(\tau) + \sum_{j=1}^J \mathbf{b}_{2j+1} B_{2j+1}(\tau) + \sum_{\nu=1}^{\bar{n}} \bar{\mathbf{d}}_{\nu} \sin(\nu\pi\tau)$$



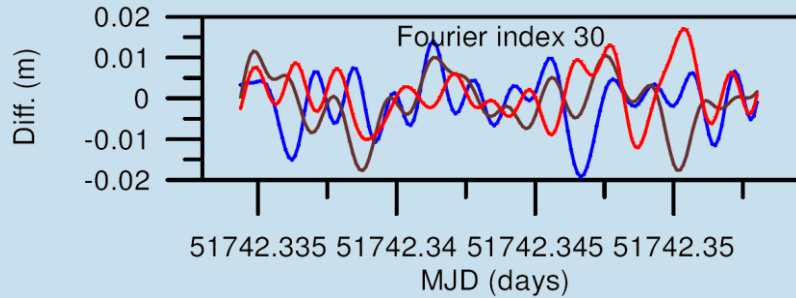
**Amplitudes**

**Remainders**

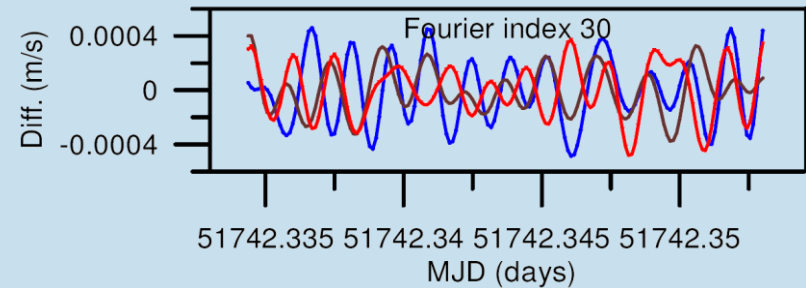
GPOD



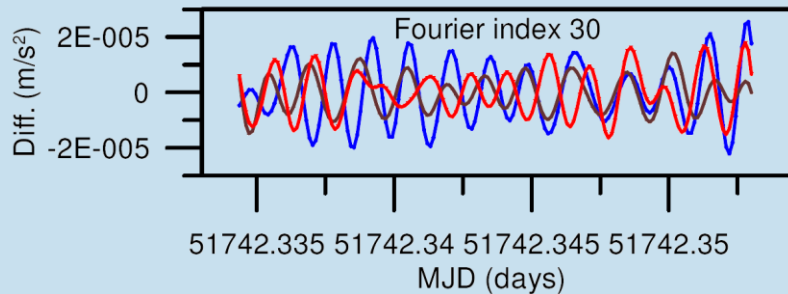
DPOD



**Position differences**



**Velocity differences**



**Acceleration differences**

RMS

index	Pos.(m)	Vel.(m/s)	Acc.(m/s <sup>2</sup> )
20	0.012644	0.000353	0.000012
30	0.010717	0.000397	0.000018
40	0.011997	0.000463	0.000025
59	0.014737	0.000941	0.000077

**Statistical values**



GPOD

KPOD

RKPOD

DPOD

Method

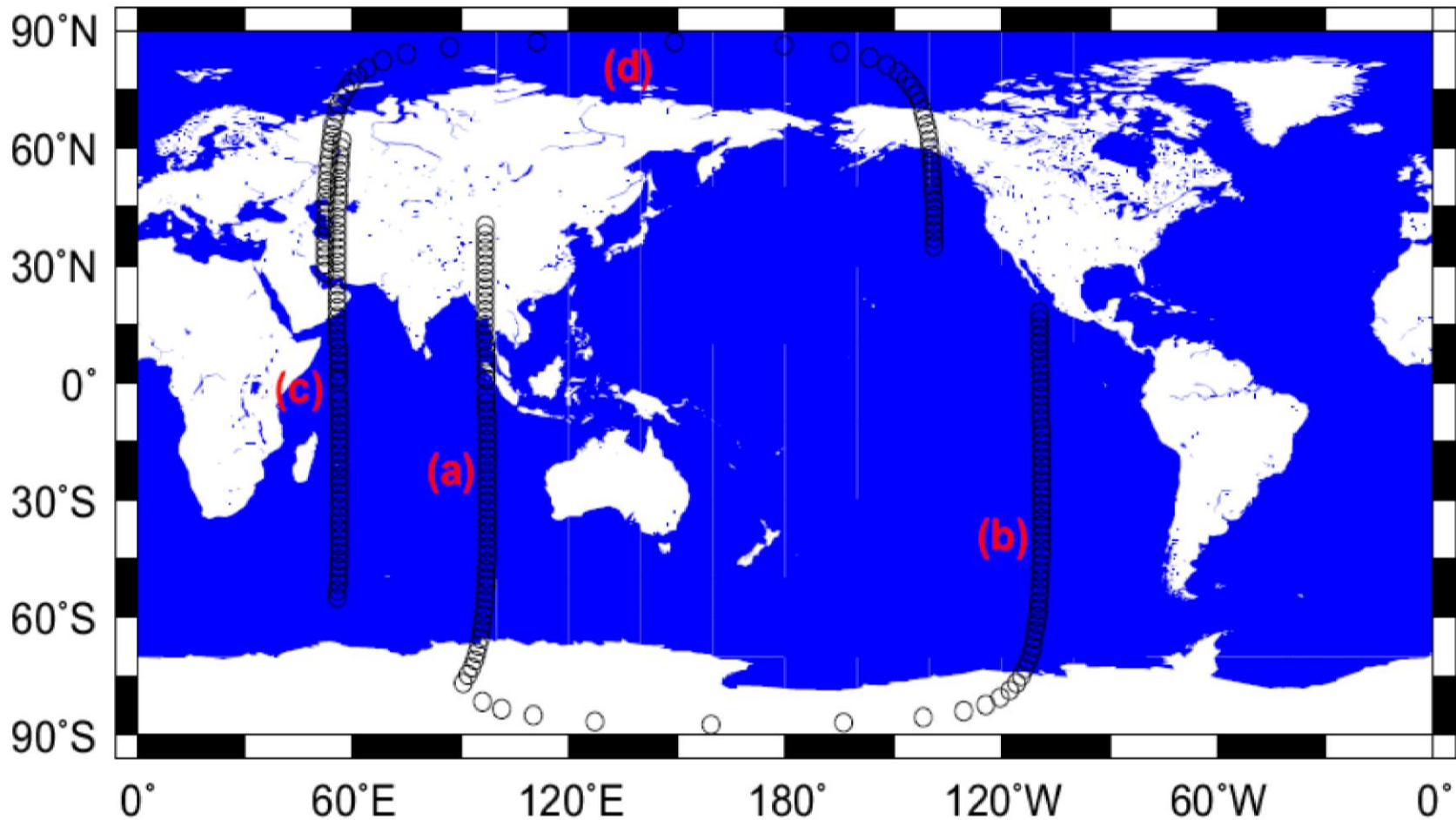
Method

Simulation

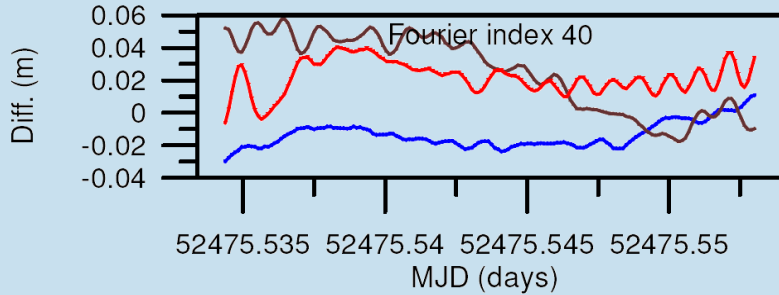
Simulation

Real data

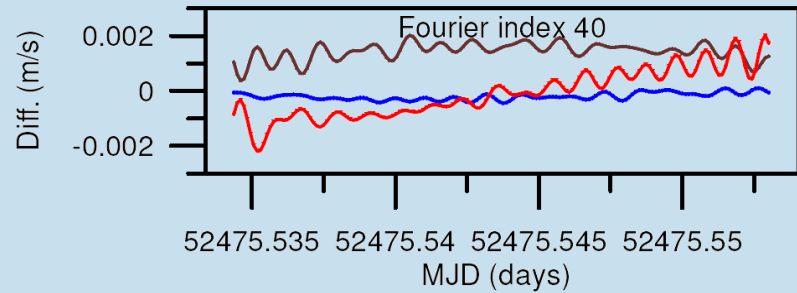
Real data



Four short arcs (30 min.) ground track of CHAMP



## IGG - GFZ positions

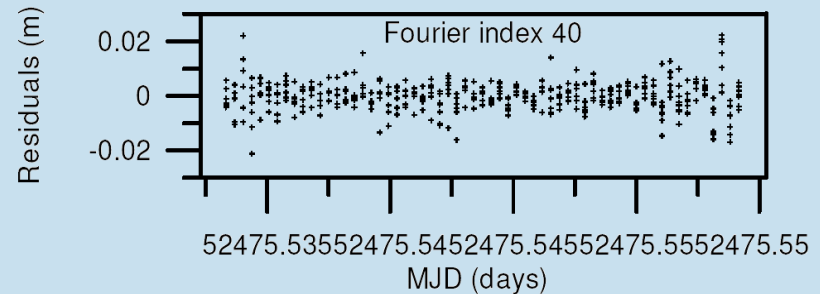


## IGG - GFZ velocities

RMS

index	Pos.( <i>m</i> )	Vel.( <i>m/s</i> )
20	0.0503	0.0019
30	0.0455	0.0018
40	0.0449	0.0017
59	0.0449	0.0017

## Statistical values



## GPS-SST residuals





GPOD

KPOD

RKPOD

DPOD

Method

Method

Simulation

Simulation

Real data

Real data

Dynamical  
info.

$$= \bar{\mathbf{r}}(\tau) + \sum_{\nu=1}^n \mathbf{d}_{\nu} \sin(\nu\pi\tau)$$

$$\tilde{\mathbf{d}}_{\nu} = -\frac{2T^2}{\nu^2\pi^2 - \mu^2} \int_{\tau'=0}^1 \sin(\nu\pi\tau) \mathbf{a}(\tau'; \mathbf{r}, \dot{\mathbf{r}}) d\tau'$$

✓ Introduction of an approximate force function  $(\tilde{\mathbf{d}}_i \cdots \tilde{\mathbf{d}}_j)$ ,  $\mathbf{C}_{(\tilde{\mathbf{d}}_i \cdots \tilde{\mathbf{d}}_j)}$

✓ Fixing only some orbit parameters  $(\tilde{\mathbf{d}}_i \cdots \tilde{\mathbf{d}}_j)$ ,  $\mathbf{C}_{(\tilde{\mathbf{d}}_i \cdots \tilde{\mathbf{d}}_j)} \rightarrow \mathbf{0}$

✓ Down- or up weighting  $\mathbf{C}_{(\tilde{\mathbf{d}}_1 \cdots \tilde{\mathbf{d}}_n)}$  in relation to  $\mathbf{C}_{(\mathbf{d}_1 \cdots \mathbf{d}_n)}$

Kinematical observation equation

$$\mathbf{I}_1 = (\mathbf{A}_1 \quad \mathbf{A}_2) \begin{pmatrix} \mathbf{x}_1 \\ \mathbf{x}_2 \end{pmatrix}, \quad \mathbf{C}_1$$

Constraints

$$\mathbf{I}_2 = (\mathbf{0} \quad \mathbf{I}) \begin{pmatrix} \mathbf{x}_1 \\ \mathbf{x}_2 \end{pmatrix}, \quad \mathbf{C}_2$$



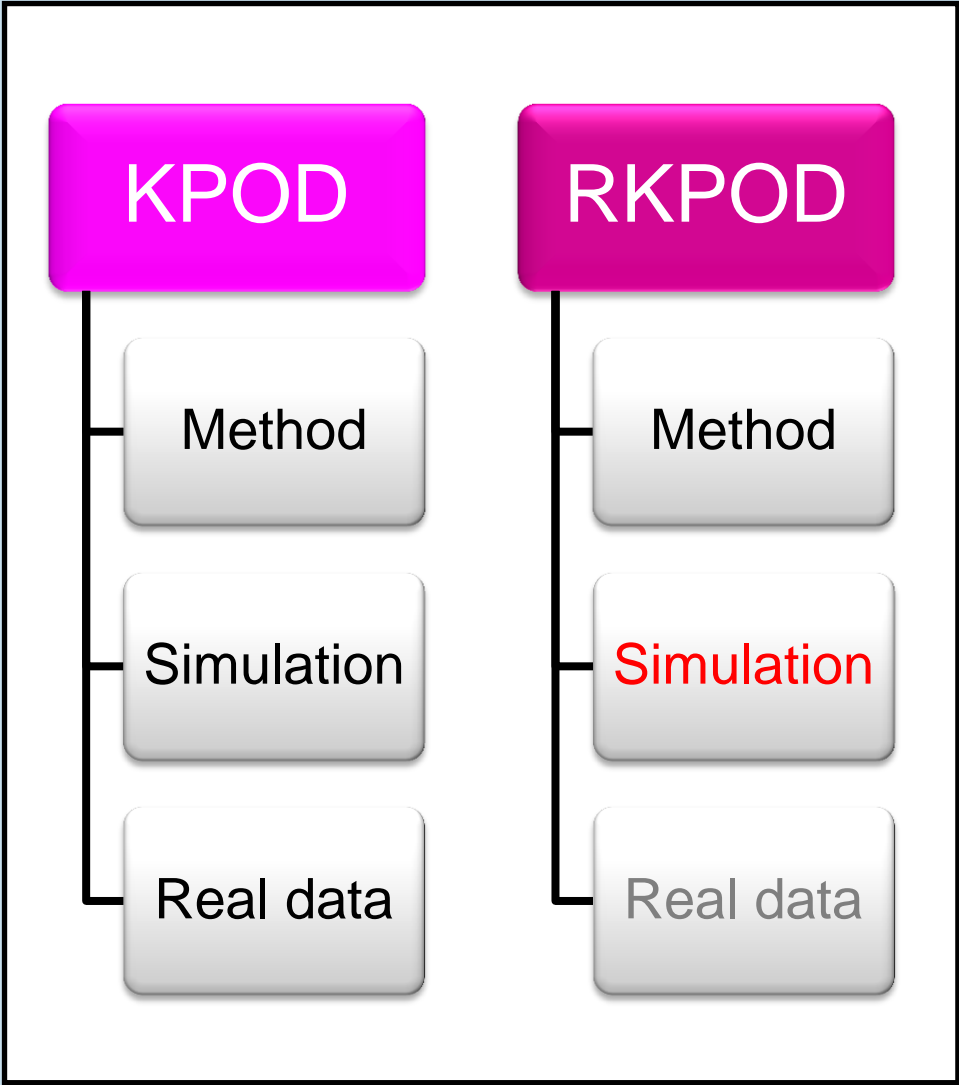
$$\begin{pmatrix} \mathbf{I}_1 \\ \mathbf{I}_2 \end{pmatrix} = \begin{pmatrix} \mathbf{A}_1 & \mathbf{A}_2 \\ \mathbf{0} & \mathbf{I} \end{pmatrix} \begin{pmatrix} \mathbf{x}_1 \\ \mathbf{x}_2 \end{pmatrix}, \quad \mathbf{C}_1 = \begin{pmatrix} \mathbf{C}_1 & \mathbf{0} \\ \mathbf{0} & \mathbf{C}_2 \end{pmatrix}$$



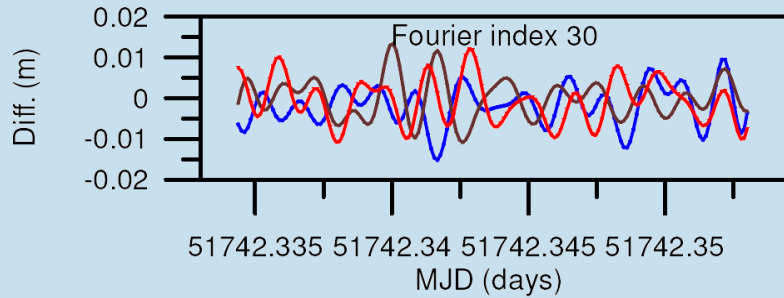
$$\begin{pmatrix} \hat{\mathbf{x}}_1 \\ \hat{\mathbf{x}}_2 \end{pmatrix} = \mathbf{N}^{-1} \begin{pmatrix} \mathbf{A}_1^T \mathbf{C}_1^{-1} \mathbf{I}_1 \\ \mathbf{A}_2^T \mathbf{C}_1^{-1} \mathbf{I}_1 + \mathbf{C}_2^{-1} \mathbf{I}_2 \end{pmatrix}, \quad \mathbf{N}^{-1} = \begin{pmatrix} \mathbf{Q}_{\hat{\mathbf{x}}_1 \hat{\mathbf{x}}_1} & \mathbf{Q}_{\hat{\mathbf{x}}_1 \hat{\mathbf{x}}_2} \\ \mathbf{Q}_{\hat{\mathbf{x}}_2 \hat{\mathbf{x}}_1} & \mathbf{Q}_{\hat{\mathbf{x}}_2 \hat{\mathbf{x}}_2} \end{pmatrix}$$



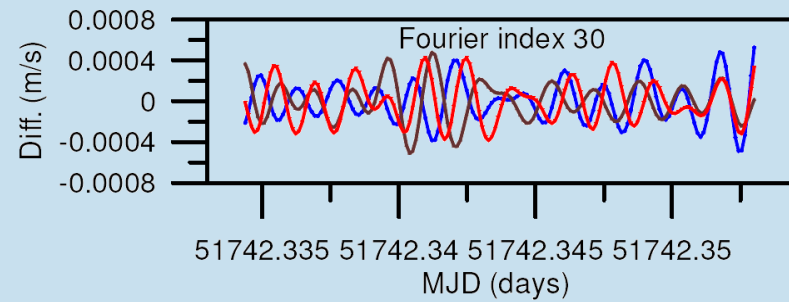
GPOD



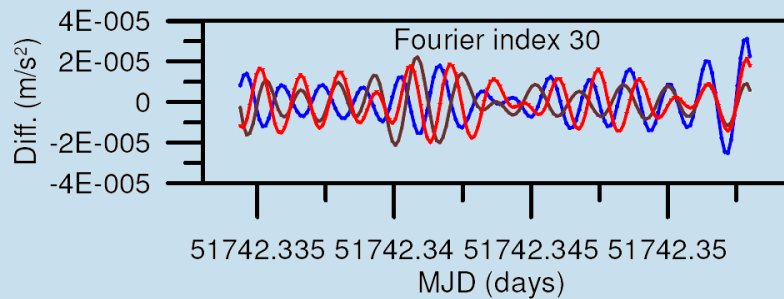
DPOD



## Position differences



## Velocity differences



## Acceleration differences

RMS

index	Pos.(m)	Vel.(m/s)	Acc.(m/s <sup>2</sup> )
20	0.012831	0.000316	0.000012
30	0.008873	0.000337	0.000016
40	0.014034	0.000402	0.000021
59	0.011553	0.000721	0.000056

## Statistical values

➤ GNSS-LEO satellites configuration and geometrical strength play an important role in POD,

➤ Kinematical POD can be used to recover the Earth's gravity field model based on the POD methods,

➤ No gravity field and no force models have been used in the Geometrical and Kinematical modes (**advantage**),

➤ The proposed kinematical orbit determination method is very flexible. A smooth transition from kinematical to reduced kinematical and finally dynamical or vice-versa is possible.

**Thank you  
for your attention**

